Fluctuation-dissipation relations for thermodynamic





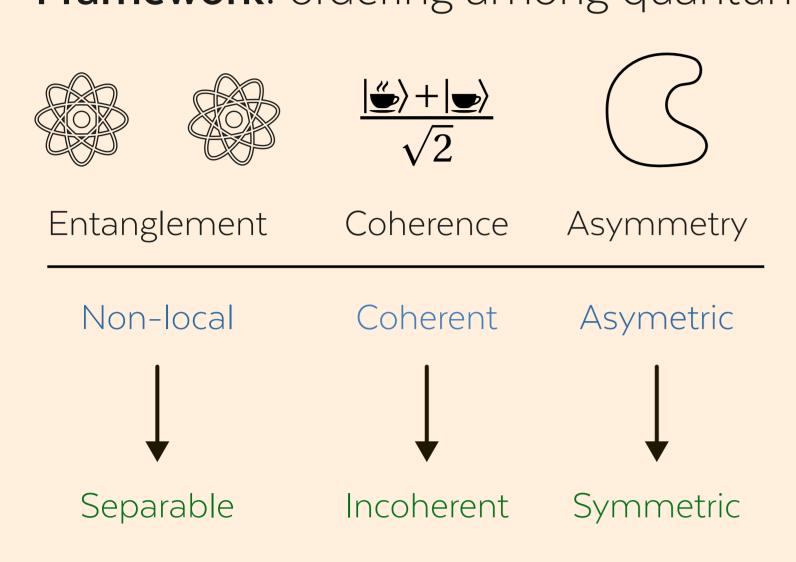
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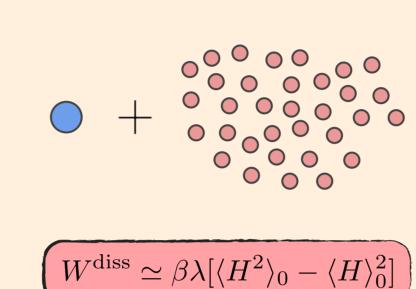
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Background

Goal: study resource dissipation and characterise optimal state transformation protocols that minimise dissipation

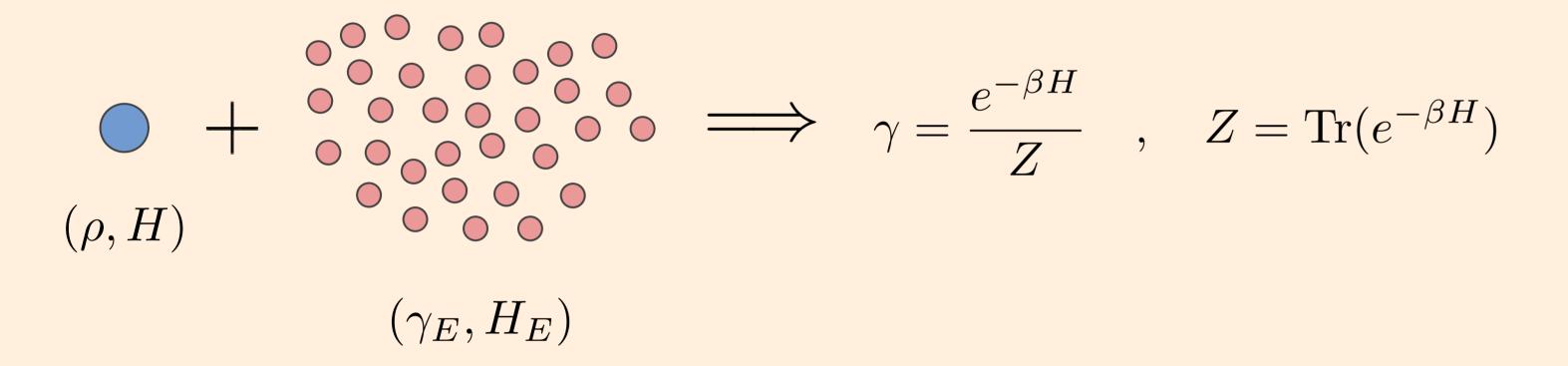
Framework: ordering among quantum states





Setting the scene

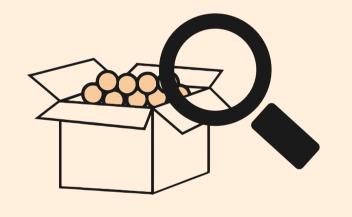
Indentifying the set of thermodynamically-free states



Thermodynamic transformations are modelled by thermal operations

$$\mathcal{E}(
ho)=\mathrm{Tr}_E(U(
ho\otimes\gamma_E)U^\dagger)$$
 with $[U,H\otimes\mathbb{1}_E+\mathbb{1}_E\otimes H_E]=0$ Energy-conserving

Thermodynamic monotone $\phi: \mathcal{S}_d \to \mathbb{R}_+ \cup \{0\}$



i.
$$\phi(\mathcal{E}(\rho)) \leq \phi(\rho)$$

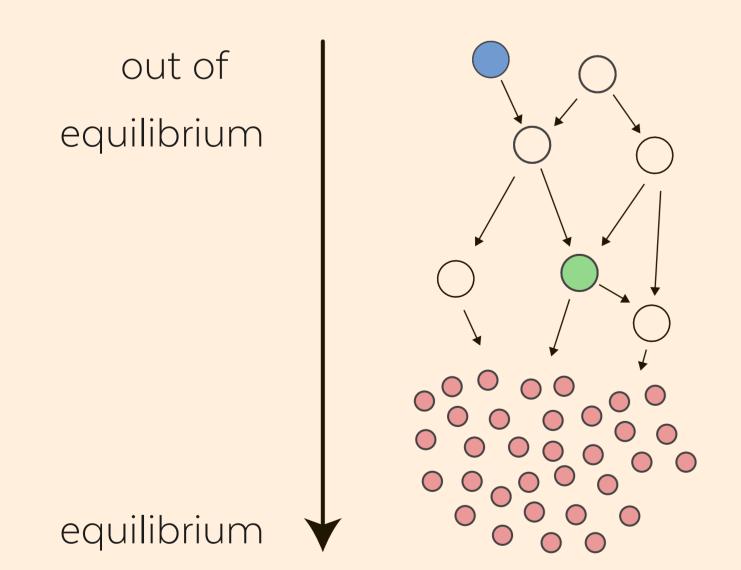
$$D(\rho || \gamma) = \text{Tr}(\rho(\log \rho - \log \gamma))$$

ii.
$$\phi(\gamma) = 0$$

Generalised free energy

Thermodynamic distillation process

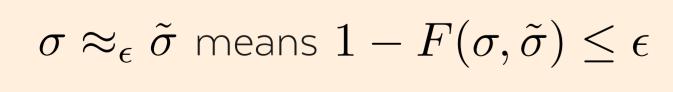
For initial state ρ , target state σ , thermal bath $\beta \Longrightarrow \mathcal{E}(\rho) = \sigma$

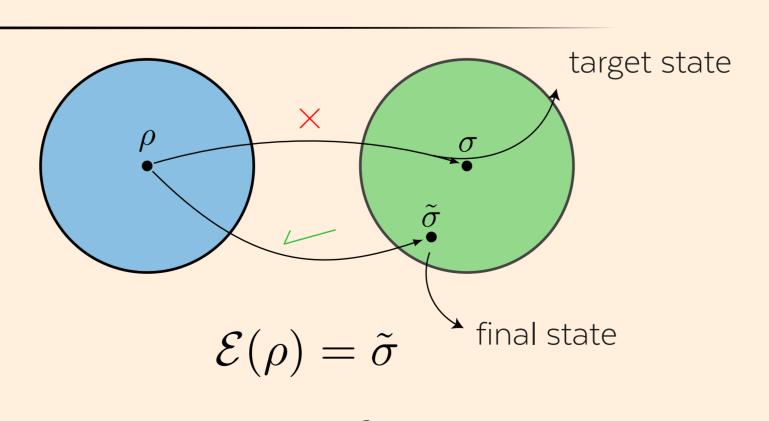


- General answer not known beyond the simplest qubit case
- For **energy-incoherent** states

$$\mathcal{E}(\rho) = \sigma : \mathbf{p} \succ^{\beta} \mathbf{q}$$

 ϵ -approximate **interconversion** problem:





 ϵ -approximate thermodynamic distillation process: $(\rho, H) \xrightarrow{\mathcal{E}} (\tilde{\rho}, \tilde{H})$

with
$$\tilde{\rho} = \bigotimes_{m=1}^N |\tilde{E}_{k_n}^{(n)}\rangle \langle \tilde{E}_{k_n}^{(n)}|$$

Eigenstate of
$$|\tilde{E}_{k_n}^{(n)}\rangle$$
 correspoding to energy $\tilde{E}_k^{(n)}$

with
$$\tilde{\rho} = \bigotimes_{m=1}^{\tilde{N}} |\tilde{E}_{k_n}^{(n)}\rangle \langle \tilde{E}_{k_n}^{(n)}|$$

$$\sum_{m=1}^{\tilde{N}} |\tilde{E}_{k_n}^{(n)}\rangle \langle \tilde{E}_{k_n}^{(n)}|$$
 Eigenstate of $|\tilde{E}_{k_n}^{(n)}\rangle$ correspoding to energy $\tilde{E}_k^{(n)}$
$$\sigma^2(F^N) := \frac{1}{\beta^2} \sum_{n=1}^N V(\rho_n^N \| \gamma_n^N)$$
 Free energy fluctuations

Results

Theorem 1. For a distillation setting with energy incoherent initial states, the transformation error of the approximate distillation process in the asymptotic limit is given by

$$\lim_{N \to \infty} \epsilon_N = \lim_{N \to \infty} \Phi \left(-\frac{\Delta F^N}{\sigma(F^N)} \right)$$

Moreover, for any N there exist an approximate distillation process with the transformation error bounded by

$$\epsilon_N \le \Phi\left(-\frac{\Delta F^N}{\sigma(F^N)}\right) + \frac{C\kappa^3(F^N)}{\sigma^3(F^N)}$$

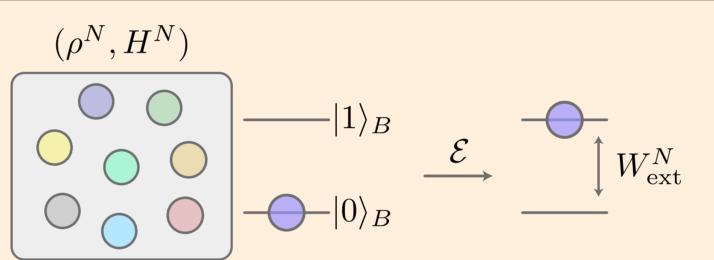
Theorem 2. For a distillation setting with N identical initial systems each in a pure state and described by the same Hamiltonian, the transformation error of the optimal approximate distillation process in the asymptotic limit is given by

$$\lim_{N \to \infty} \epsilon_N = \lim_{N \to \infty} \Phi\left(-\frac{\Delta F^N}{\sigma(F^N)}\right)$$

The amount of **free energy dissipated** in both settings satisfies:

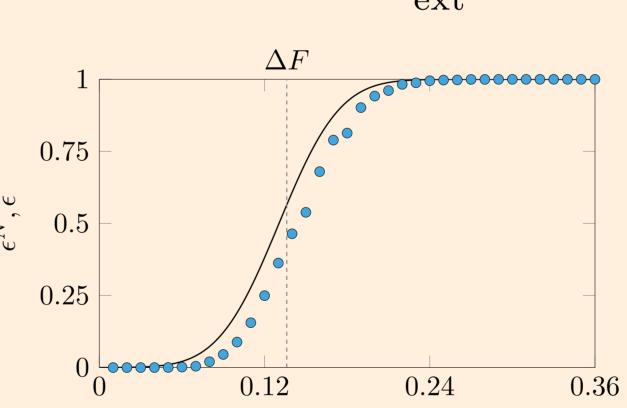
$$F_{\rm diss}^{\rm tot} = a(\epsilon)\,\sigma^{\rm tot}(F) \quad \text{with } a(\epsilon) = -\Phi^{-1}(\epsilon)(1-\epsilon) + \frac{\exp\left(\frac{-(\Phi^{-1}(\epsilon))^2}{2}\right)}{\sqrt{2\pi}}$$

Aplications



Optimal work extraction

■ Theorem 1 tell us that the optimal transformation error for extracting the amount of work $W_{
m ext}^N$ is:



 $w_{\mathrm{ext}}^N, w_{\mathrm{ext}}$

Initial system is composed of 100 two-level subsystems. The first 59 is prepared in the state p = (0.9, 0.1) with thermal state (0.6,0.4) and remaning is prepared in a state (0.7, 0.3) and thermal state (0.75, 0.25).

 $W_{\mathrm{ext}}^N \simeq F^N + \sigma(F^N)\Phi^{-1}(\epsilon)$

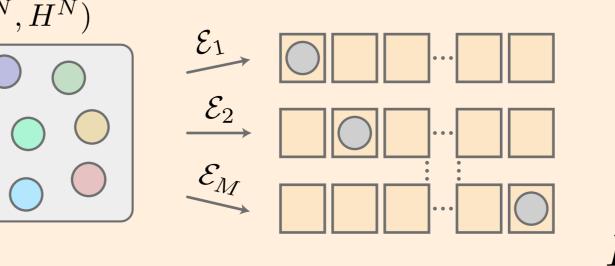
■ From Theorem 2, the optimal amount of work extracted from N pure quantum systems up to second-order is given by

$$W_{\rm ext} \simeq N \left(\langle H \rangle_{\psi} + \frac{\log Z}{\beta} + \frac{\langle H^2 \rangle_{\psi} - \langle H \rangle_{\psi}^2}{\sqrt{N}} \Phi^{-1}(\epsilon) \right)$$

Optimal information erasure

applying Theorem 1 yields: $\| \Psi_{\mathrm{cost}}^{-|1\rangle_B} \| \mathcal{E} \| \| \Psi_{\mathrm{cost}}^{N} \| W_{\mathrm{cost}}^{N} \| W_{\mathrm{cost}}^{N} \| \mathcal{E} \| \| \mathcal{E} \| \| \Psi_{\mathrm{cost}}^{N} \| W_{\mathrm{cost}}^{N} \| \mathcal{E} \| \| \mathcal{E} \| \| \mathcal{E} \| \| \Psi_{\mathrm{cost}}^{N} \| W_{\mathrm{cost}}^{N} \| \mathcal{E} \| \mathcal{E} \| \mathcal{E} \| \| \mathcal{E}$

Optimal thermodynamically-free communication rate



The optimal number of messages that can be encoded into in a thermodynamically-free

$$R(\rho^{\otimes N}, \epsilon_{\mathrm{d}}) \simeq D(\rho \| \gamma) + \frac{\sqrt{V(\rho \| \gamma)}}{\sqrt{N}} \Phi^{-1}(\epsilon_{\mathrm{d}})$$

In this result is valid for either a pure or incoherent state!