## Stochastic Path

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Abstract and outline

We consider the evolution of a harmonic oscillator in Gaussian states undergoing simultaneous weak conti and momentum measurements. The expectation value and momentum show stochastic evolution while the matrix elements show deterministic evolution and co steady states. We apply the Chantasri-Dressel-Jorda integral formalism to express the probability density readouts as a path integral of the exponential of a st Extremization of the action gives us the most-likely the analytical solutions for these most-likely paths for state values of the covariance matrix elements. We  $\epsilon$ energetics of the measurement process and character probability densities starting from an initial state. results using simulations.

### System description

- System Hamiltonian is  $\hat{H} = \frac{1}{2}(\hat{X}^2 + \hat{P}^2).$
- Weak continuous measurements (in d au intervals) of Gaussian operators

$$\hat{M}_X(r_1) = \left(\frac{d\tau}{2\pi\mathcal{T}_1}\right)^{\frac{1}{4}} \exp\left[-\frac{d\tau}{4\mathcal{T}_1}(r_1\mathbb{1} - \hat{X})^2\right] \qquad \hat{M}_P(r_2) = \left(\frac{d\tau}{2\pi\mathcal{T}_2}\right)^{\frac{1}{4}} \exp\left[-\frac{d\tau}{4\mathcal{T}_2}(r_1\mathbb{1} - \hat{X})^2\right]$$

- $r_1, r_2$  are the readouts of position and momentum measuren
  - $\mathcal{T}_1$  and  $\mathcal{T}_2 \ (\gg d au)$ scales of the X and
  - The unitary dynamical dynami operators preserve the state.
  - We define  $\langle \hat{X} \rangle$ ,  $\langle$

 $2\mathrm{Cov}(\hat{X},\hat{R})$  $2\operatorname{Var}(X)$ • Covariance matrix  $\Gamma =$  $\left\langle \hat{X}\hat{P} + \hat{P}\hat{X} \right\rangle - 2\left\langle \hat{X} \right\rangle \left\langle \hat{P} \right\rangle \quad 2 \text{Var}(\hat{P})$ For pure states det  $\Gamma = 1$ 

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State evolution

• Sample trajectories from the same initial state are shown for  $\mathcal{T}_1 = \mathcal{T}_2 = 1$ and  $\tau_f = 5.00$ .

Harmonic Oscillator

- In this case,  $\Gamma$  converges to the  $\mathbb{1}_2$  at  $\tau \to \infty$ . For simplicity we assume  $\Gamma = \mathbb{1}_2.$
- The quadratures  $(\mathbf{q})$  show stochasticity due to measurement back action.
- Spiral trajectories in phase space signify energy gain/loss due to measurements.



Integral A	Analysis of the Continuously Monitored S	imple Har
NV UCA 2 Conton for Cohomor	Tathagata Karmakar <sup>1,2,3</sup> , Philippe Lewalle <sup>1,2,4</sup> , Andrew N. Jordan <sup>1,2,3</sup>	University Onenana CA UCA
NY, USA, - Center for Conerer	Chapter i Dregel Ierden stochester, NY, USA, Institute for Quantum Studies, Chapman (	University, Orange, CA, USA,
1	The CDI formalism lots us calculate the probability density associated with a	
n general	• The ODJ formatism lets us calculate the probability density associated with a	
inuous position	sequence of readouts $\{T_k\}$ .	• The conjugate mor
ues of position	• $q(0) = q_i \rightarrow q(\tau_f) = q_f$ through intermediate states $\{q_k\}$ and readouts $\{r\}$ .	while the energy co
onverge to their	• The CDJ formalism tells us $\int \int \int \tau_f$	corresponds to $\boldsymbol{p} =$
an stochastic nath	$\mathcal{P}(\{\boldsymbol{q}_k\}, \{\boldsymbol{r}_k\}) = \left  \mathcal{D}\boldsymbol{p}  e^{\mathcal{S}[\boldsymbol{p}, \boldsymbol{q}, \boldsymbol{r}]} = \left  \mathcal{D}\boldsymbol{p}  \exp \left  \int^{T} d\tau (-\boldsymbol{p} \cdot \dot{\boldsymbol{q}} + \mathcal{H}(\boldsymbol{p}, \boldsymbol{q}, \boldsymbol{r})) \right .$	• The expectation va
v of a sequence of	$J \qquad J \qquad$	mechanical energy
tochastic action.	p is momentum	$E_{\text{res}}(\pi) = \frac{1}{1} + \frac{1}{5} \left( \frac{1}{5} \right)$
paths. We find	<b>S</b> and <b>21</b> and the stackestic setion and stackestic Hereiter mean estimate	$E_M(7) = \frac{1}{2} + \frac{1}{2} \{ \{ \mathcal{Q}_{1j} \} \}$
or the steady	• $\mathcal{S}$ and $\mathcal{H}$ are the stochastic action and stochastic Hamiltonian respectively.	$\left(\mathcal{Q}_{2f}+\frac{\tau}{-}(q_{2f}-\mathcal{Q}_{2f})\right)$
also analyze the	• $\delta \delta = 0$ gives the most-likely readouts. The corresponding trajectory is the	$\mathcal{O}_{1f} = a_{1i}\cos \tau_f + a_{2i}\sin \theta$
rize final state	most-likely path or the optimal path.	$O_{2i} = -a_1 \cdot \sin \tau i + a_2$
Ve confirm our	• For the steady state of the covariance matrix elements, the most-likely paths can	$\mathbf{z}_{2f} - \mathbf{q}_{1i} \sin \mathbf{r}_{f} + \mathbf{q}_{2}$ final state along
	be solved for analytically.	opporting optimal
	Optimal paths	conserving optimal
	• The energy conserving optimal path (yellow solid) is circular in phase space, as	H A A A A A A A A A A A A A A A A A A A
n type with Kraus	the distance from the origin corresponds to the expectation value of the	Using the CDJ path :
	mechanical energy. The energy decreasing (red dash-dotted) and energy	
$\left  -\frac{d\tau}{r_2} (r_2 \mathbb{1} - \hat{P})^2 \right .$	increasing (red dashed) optimal paths are spirals.	
$\begin{bmatrix} 4\mathcal{T}_2 \\ \end{pmatrix}$	$ au_1 = 1.0, \  au_2 = 1.0$	
nents.	-0.30	
		0.06
) are collapse time	5 $-\frac{1}{2}$ $i(\tau = 0.25)$	
nd $P$ measurements.		$\widehat{\mathbf{p}}^{0.04}$
mics and the Kraus		
e the Gaussianity of		-0.02
J	$+$ $+$ $\tau = 3.0$ $+$ $-0.10$	
$\langle \hat{P} \rangle$ ) <sup>T</sup> = $(a_1 \ a_2)^{T} = \boldsymbol{a}$	$\tau = 2.0$	-5
$\langle 1 \rangle \langle (q_1, q_2) \rangle q_1$		$q_2( au$
<b>\</b>		
() 1.	$0$ $2$ $4_4$	
		• Using Chantasri-D
	$\Delta E_M < 0  2 -$	a continuous varial
	$- \Delta E_M > 0$	statistics of a conti
	$\Delta E_M = 0 \qquad 0$	- Our optimal path
	-5 0 $5$	foodbook control of
	$q_1$	Orman Control al
	A compute stackartic traineterry (massa) is also shown	• Our work also prov
	• A sample stochastic trajectory (green) is also shown.	novel quantum tec
	• Averages of simulated clustered trajectories (denoted as + ) for the same	
	boundary conditions show very good agreement with the analytical results.	• I. Karmakar, P. L
	• Histograms show the diffusion of trajectories at three different times. The	• A. Chantasri, J. D
	colorbar denotes the probability density of trajectories.	• K. Jacobs and D. A
	• The energy conserving trajectory is the the globally most-likely path.	• X. B. Wang, T. Hi
	DOCUTED JOHN	(2007).
$q_1( au)$	FOUNDATION FOUNDATION	• E. Arthurs and J.



# monic Oscillator

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Final state probability densities integral formalism we also get  $P(\boldsymbol{q}_f | \boldsymbol{q}_i)$ .



• For 
$$\mathcal{T}_1 = \mathcal{T}_2 = \mathcal{T}$$
 and  $\Gamma = \mathbb{1}_2$ ,  
 $P(\boldsymbol{q}_f | \boldsymbol{q}_i) =$   
 $\frac{2\mathcal{T}}{\pi \tau_f} \exp\left\{-\frac{2\mathcal{T}}{\tau_f}(q_{1f} - \mathcal{Q}_{1f})^2\right\}$   
 $\times \exp\left\{-\frac{2\mathcal{T}}{\tau_f}(q_{2f} - \mathcal{Q}_{2f})^2\right\}$ 

• Analytical results (orange surface) show good agreement with simulated results (black wireframe).

### Summary and discussions

ressel-Jordan stochastic path integral formalism, first time for ble system, we completely characterize the measurement inuously monitored general Gaussian state harmonic oscillator. description and analysis of the energetics can be useful for e.g. nd cooling of resonators.

vides a new way to connect optomechanical elements with chnologies like quantum measurement engines and refrigerators.

#### References

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